

Three-transition cascade erbium laser at 1.7, 2.7, and 1.6 μm

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We report on an upconversion cascade laser in an erbium-doped ZBLAN fiber emitting simultaneously on the three transitions ${}^4S_{3/2} \rightarrow {}^4I_{9/2}$ at 1.7 μm , ${}^4I_{11/2} \rightarrow {}^4I_{13/2}$ at 2.7 μm , and ${}^4I_{13/2} \rightarrow {}^4I_{15/2}$ at 1.6 μm . At moderate pump powers, the laser transition at 1.6 μm supports 2.7- μm lasing and permits a slope efficiency at 2.7 μm of 15% versus launched pump power. Above the threshold of upconversion lasing at 1.7 μm , the slope efficiency at 2.7 μm increases to 25.4%. Taking pump excited-state absorption into account, this value represents more than 90% of the theoretical slope efficiency. A transversely single-mode output power of 99 mW is achieved at 2.7 μm . © 1997 Optical Society of America

The increasing interest in 2.7–2.8- μm lasers is evoked by applications in surgery.^{1,2} Erbium-doped fluoride fibers are promising candidates for the construction of compact and efficient all-solid-state laser sources emitting between 2.7 and 2.8 μm . In diode-pumped double-clad fiber systems, generally a high-brightness transversely single-mode laser output can be achieved.³ The resulting high output intensity is favorable for the ablation process in surgery.

In earlier research, a saturation of the output power of the fiber laser owing to competitive lasing at 850 nm was observed.⁴ The saturation was recently overcome in an upconversion cascade lasing regime⁵ at 1.7 and 2.7 μm by energy recycling into the 2.7- μm laser process. Stringent demands on the pump intensity, however, have prevented this system from double-clad pumping with low-brightness diode lasers. Other possibilities for overcoming the saturation effect are the quenching of the lower laser level by colasing at 1.6 μm (Ref. 6) or a combination of this approach with the upconversion cascade regime of Ref. 5, which results in lasing on three transitions.

In this Letter we report on a three-transition cascade laser emitting simultaneously on the transitions ${}^4S_{3/2} \rightarrow {}^4I_{9/2}$ at 1.7 μm , ${}^4I_{11/2} \rightarrow {}^4I_{13/2}$ at 2.7 μm , and ${}^4I_{13/2} \rightarrow {}^4I_{15/2}$ at 1.6 μm (see Fig. 1). The quenching of the lower laser level of the 2.7- μm transition is achieved by a combination of cascade lasing at 1.6 μm and pump excited-state absorption⁷ (ESA). The energy accumulated in the ${}^4S_{3/2}$ level is recycled into the laser process at 2.7 μm by cascade lasing at 1.7 μm . A slope efficiency of 25.4% is obtained at 2.7 μm , which is, to our knowledge, the highest value reported so far for an erbium 2.7- μm fiber laser. Differences between this three-transition cascade regime and the performance of a two-transition cascade regime at 1.7 and 2.7 μm are investigated.

A fluorozirconate fiber (labeled fiber A; see Table 1) is pumped by a Ti:sapphire laser at 792 nm. The pump beam is chopped with a frequency of 10 Hz and a duty cycle of 50%. Approximately 55% of the incident

pump power is launched into the fiber. This value includes a transmission of the incoupling optics of 94%, a transmission of the incoupling mirror of 84%, and an estimated coupling efficiency into the fiber of 70%. The fiber length of 1.1 m minimizes reabsorption losses on the ground-state transition and ensures high pump absorption. More than 96% of the launched pump power is absorbed in the fiber through ground-state absorption (GSA) and ESA (see Fig. 1).

The mirrors are butt coupled to the fiber ends. Mirror transmissions are 88% at 850 nm, 1% at 1.6 μm , 1% at 1.7 μm , and 68% at 2.7 μm . High transmission at 850 nm increases the threshold of the high-gain 850-nm laser. Both the 1.6- and the 1.7- μm mirror transmissions are designed for high intracavity power rather than high output power to support lasing at 2.7 μm . The mirrors could not be optimized for the

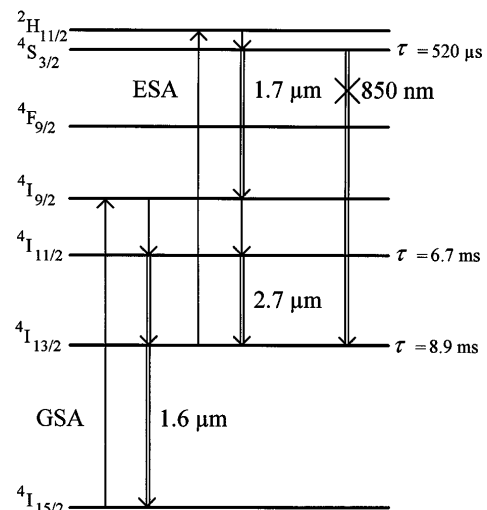


Fig. 1. Energy-level scheme of Er^{3+} in ZBLAN fiber. The system is pumped by GSA and ESA at 792 nm. Laser emission is obtained on the three transitions ${}^4S_{3/2} \rightarrow {}^4I_{9/2}$ at 1.7 μm , ${}^4I_{11/2} \rightarrow {}^4I_{13/2}$ at 2.7 μm , and ${}^4I_{13/2} \rightarrow {}^4I_{15/2}$ at 1.6 μm . Competitive lasing at 850 nm is suppressed.

Table 1. Parameters of the Investigated ZBLAN Fibers^a

Parameters	Fiber A	Fiber B
Length (m)	1.1	4.8
Core diameter (μm)	6.0	6.5
N.A.	0.4	0.16
Er ³⁺ concentration (parts in 10 ⁶ mol.)	3000	1000
Slope efficiency 1	15%	7%
1.7- μm threshold (mW)	304	120
Slope efficiency 2	25.4%	22.6%

^aFibers are from Le Verre Fluoré. Lasing at 2.7 and 1.7 μm is obtained in both fibers; lasing at 1.6 μm is only in fiber A. Comparison of the slope efficiencies obtained at 2.7 μm : slope efficiency 1 is obtained below the threshold of upconversion lasing at 1.7 μm (i.e., lasing at 2.7 and 1.6 μm in fiber A, lasing at 2.7 μm in fiber B); slope efficiency 2 is obtained with 1.7- μm lasing.

highest 2.7- μm efficiency because the requirements at 792 nm, 850 nm, 1.6 μm , and 1.7 μm had to be matched.

The input-output characteristics of the laser are shown in Fig. 2. The threshold of 2.7- and 1.6- μm lasing is at 60-mW launched pump power. The slope efficiency at 2.7 μm is 15% in the two-transition cascade regime with 1.6- μm colasing but increases to 25.4% at the onset of 1.7- μm lasing and energy recycling into the 2.7- μm laser process above 304-mW launched pump power. The Stokes limit of the system for the pump wavelength of 792 nm is 29.3%. The theoretical limit of the slope efficiency, however, is slightly smaller than the Stokes limit, because there is a small loss owing to pump ESA from the ⁴I_{11/2} level (not indicated in Fig. 1). Pump photons absorbed on this transition are not converted into laser photons at 2.7 μm , which decreases the quantum efficiency and reduces the limit of the slope efficiency to 27.5%.⁸ The experimentally obtained value of 25.4% represents 92% of this maximum possible slope efficiency.

A transversely single-mode output power of 99 mW at 2.7 μm is achieved with 553-mW launched pump power. Above this pump-power level, an increase of the spot size of the Ti:sapphire laser as well as damage of the self-fabricated mirror coating is observed. The laser lines at 1.6 and 1.7 μm have output powers of approximately 1 mW because of the high mirror reflectance at these wavelengths. The three laser lines (see Fig. 3) are centered at 1.602, 1.716, and 2.702 μm .

These results are compared with results of a two-transition cascade regime at 1.7 and 2.7 μm , which are obtained with a different fiber (labeled fiber B; see Table 1). Several remarks are necessary: (i) The same resonator mirrors are used for both experiments. (ii) The product of the fiber length and the dopant concentration is smaller by a factor of 1.45 for fiber A; this leads to a reduced ground-state reabsorption that allows for lasing at 1.6 μm , which is not achieved with fiber B. (iii) The fraction of launched pump power absorbed in the fiber is larger than 96% in both experiments. (iv) Fiber A has improved parameters of core diameter and N.A. for better guiding of the 2.7- μm laser line.

The improved fiber parameters as well as colasing at 1.6 μm increase the 2.7- μm slope efficiency at lower pump power from 7% (fiber B) to 15% (fiber A). The 1.6- μm laser repopulates the ground state and ensures that a high fraction of pump power is absorbed on the GSA transition that feeds the 2.7- μm laser. The threshold of 1.7- μm lasing and energy recycling into the 2.7- μm laser process, however, increases from 120 mW (fiber B) to 304 mW (fiber A). This is caused by the additional depletion of the ⁴I_{13/2} level through 1.6- μm lasing, which weakens the pump ESA from this level and decreases the excitation of the ⁴S_{3/2} level from which the 1.7- μm laser originates.⁹

In the upconversion cascade regime with 1.7- μm lasing, the slope efficiency at 2.7 μm is increased from 22.6% (fiber B) to 25.4% (fiber A) because of the improved fiber parameters. Additional lasing on the third transition at 1.6 μm transfers population from the ⁴I_{13/2} level to the ⁴I_{15/2} level. Whether this influences the laser performance in the upconversion cascade regime depends on the fraction of launched pump power that is absorbed within the fiber. The ESA transition from the ⁴I_{13/2} level has a higher absorption cross section than the GSA transition.⁴ If not all the pump power is absorbed, it is potentially better to maintain the population in the ⁴I_{13/2} level rather than to transfer it to the ground state, because then the increase in ESA exceeds the decrease in GSA.⁸ In the present cases, most of the launched

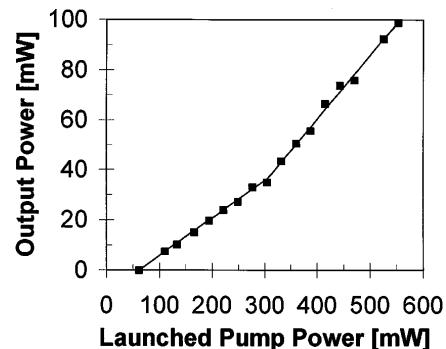


Fig. 2. Input-output curve at 2.7 μm obtained with fiber A operating in two different cascade regimes. Below 304-mW launched pump power: emission at 2.7 and 1.6 μm , slope efficiency $\eta = 15\%$. Above 304 mW: emission at 1.7, 2.7, and 1.6 μm , $\eta = 25.4\%$.

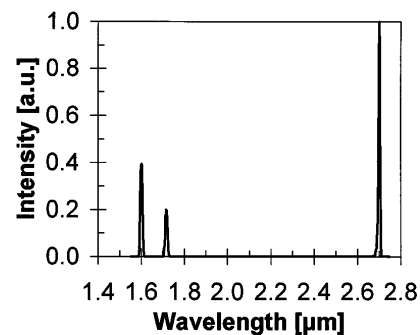


Fig. 3. Spectral behavior of the emission at 1.602, 1.716, and 2.702 μm . The intensities of the 1.6- and 1.7- μm lines are magnified by a factor of 20.

pump power is absorbed within the fiber. It is of no importance whether a pump photon is absorbed through GSA or ESA from the lower laser level. The pump photon is converted into a laser photon at $2.7 \mu\text{m}$ in both cases (see Fig. 1). Thus the $1.6\text{-}\mu\text{m}$ laser does not influence the slope efficiency in the upconversion cascade regime.

In conclusion, we have demonstrated an erbium ZBLAN fiber cascade laser operating on the three transitions $^4S_{3/2} \rightarrow ^4I_{9/2}$ at $1.7 \mu\text{m}$, $^4I_{11/2} \rightarrow ^4I_{13/2}$ at $2.7 \mu\text{m}$, and $^4I_{13/2} \rightarrow ^4I_{15/2}$ at $1.6 \mu\text{m}$. A slope efficiency of 25.4% is achieved at $2.7 \mu\text{m}$. This represents more than 90% of the possible slope efficiency at the pump wavelength of 792 nm. A transversely single-mode output power of 99 mW is obtained at $2.7 \mu\text{m}$. Both the 1.7 and the $1.6\text{-}\mu\text{m}$ transitions improve the performance of the laser operating solely at $2.7 \mu\text{m}$. Comparison with a two-transition cascade laser regime at 1.7 and $2.7 \mu\text{m}$ shows that colasing at $1.6 \mu\text{m}$ is advantageous in the low-power region below the $1.7\text{-}\mu\text{m}$ threshold but that the threshold for energy recycling into the $2.7\text{-}\mu\text{m}$ laser process through $1.7\text{-}\mu\text{m}$ lasing increases. Above the threshold of $1.7\text{-}\mu\text{m}$ lasing, the $1.6\text{-}\mu\text{m}$ laser does not influence the system. The slope efficiency at $2.7 \mu\text{m}$ in the high-power region is enhanced by improved fiber parameters.

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