

A THERMO-PNEUMATIC ACTUATION PRINCIPLE FOR A MICROMINIATURE PUMP AND OTHER MICROMECHANICAL DEVICES

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Abstract

A new type of actuator for a microminiature pump is presented. The pressure of air in a cavity, raised by resistive heating, deflects a thin silicon membrane. The dynamics of membrane deflection is studied experimentally, the results being in excellent agreement with simulation. We conclude that the device is suitable as an actuator in a micro-miniature pump.

Introduction

Since 1983, research on microminiature pumps has been carried out at the University of Twente. Several prototypes were realized in silicon, all piezoelectrically driven [1]. Major drawbacks of piezoelectric actuation are the necessity for high electrical voltages (100 V or more) and the unavailability of piezoelectric thin-film materials suitable for operation at low frequencies (10 Hz or less). Considering various actuation principles [2], it was concluded that thermo-pneumatic actuation, as we call it, might be a good alternative.

In general, we define a thermo-pneumatic actuator as a device consisting of a sealed cavity filled with a thermally expandable medium that can be heated or cooled down, resulting in a pressure change in the cavity. This induced pressure change is used to actuate a flexible or movable part, like a membrane or piston. Gas or a gas/liquid system can be used as the expanding medium. A temperature change in the cavity can be achieved by resistive heating, the use of a Peltier element or a laser. This actuation principle has been used in a microminiature valve, using a gas/liquid system and resistive heating [3]. This paper reports on the realization, behaviour and simulation of a thermo-pneumatic actuator designed to drive a microminiature pump.

Figure 1 shows a sketch of the actuator, comprising a cavity filled with air and a built-in aluminium meander that serves as a heater resistor. The resistor is supported by a thin silicon sheet suspended by four small silicon beams, which serve as a restriction to the heat flow. Aluminium current leads

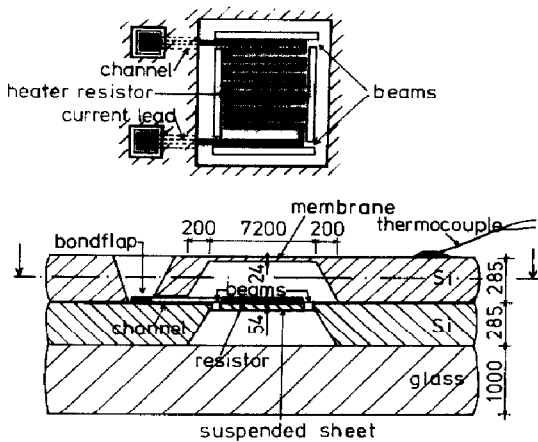


Fig. 1. Sketch of the thermo-pneumatic actuator (dimensions in μm).

connecting the meander to the bond flaps run through narrow channels, which form a restriction to the gas flow. This actuator is intended to be mounted on a pump, such that the flexible membrane can displace liquid present in a pump chamber.

Simulation model

Bond graph techniques [4] and TUTSIM [5] are used for physical modelling and simulation of the dynamic behaviour of the actuator. The model comprises the thermal behaviour of the actuator, the thermodynamics of the gas, the fluid dynamics of the channels and the mechanics of the flexible membrane. A detailed description of the simulation model will be reported elsewhere [2].

Experiments

Starting materials for the actuator are two silicon (100) 2 inch wafers, polished on both sides, and a Duran borosilicate glass wafer. The silicon wafers are shaped by wet chemical etching in a KOH water solution using standard photolithography for pattern definition. The glass wafer is cut out off a plate and polished (surface roughness $<0.05 \mu\text{m}$). The wafers are attached to one another by means of anodic bonding [6, 7]. To achieve a bond between the two silicon wafers, an intermediate layer of silicon oxide and sputtered borosilicate glass is required. A thermocouple is attached to the upper silicon wafer to monitor the temperature of the actuator as a whole.

Deflection of the flexible membrane as a function of power dissipated in the resistor is measured dynamically using a DEKTAK 3030 surface profiler scanning over $50 \mu\text{m}$ during 50 s with the scanning needle approx-

imately at the centre of the membrane. The scan length being small compared with the membrane width, the needle remains virtually in the same position. Thus, deflection at the centre of the membrane, w_0 , versus time is plotted directly.

Results and discussion

Figure 2 shows w_0 as a function of time, t , for a number of d.c. voltages, U , applied to the resistor. The voltage is switched on at time $t = 4$ s. A fast increase with a relaxation time τ_t is followed by a slower decrease to zero deflection, with a relaxation time τ_p . The first fast increase is caused by a temperature rise in the cavity and a related pressure increase. The relaxation time for warming up the cavity is mainly determined by the thermal conductance, G_{ta} , of the air in the cavity, the thermal conductance, G_{tb} , of the small suspending beams, and by the total heat capacity of the cavity. Since the heat capacity of the air in the cavity is negligible, this total heat capacity approximately equals the heat capacity, C_t , of the suspended thin silicon sheet supporting the resistor. Maximum temperature and pressure build-up and related deflection w_0 are determined by input power and thermal conductances G_{ta} and G_{tb} .

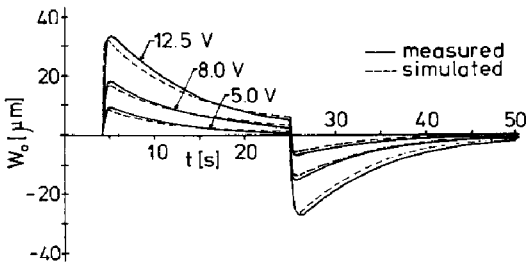


Fig. 2. Deflection w_0 vs. time for a number of applied voltages.

The subsequent slower decrease of w_0 is due to a flow of gas out of the cavity and a related pressure decrease. τ_p should be determined by the volume of the cavity, V_p , and the flow conductances, G_{pc} , of the channels. However, the calculated value of G_{pc} is much too low to account for τ_p . Evidently there is some leak in our test device, with a flow conductance G_{pl} , which was confirmed by the observation of a partial failure during the anodic bonding of the two silicon wafers.

When the d.c. voltage is switched off at time $t = 25$ s, the opposite behaviour is observed. The cavity cools down, causing a pressure drop in the cavity and a related fast decrease of w_0 , with a relaxation time τ_t' approximately equalling τ_t . The following slower increase to zero deflection is due to a flow of gas into the cavity, with a relaxation time τ_p' determined by V_p and G_{pl} , and therefore approximately equalling τ_p . The dashed lines in Fig. 2 represent results of simulations. Table 1 summarizes some parameter values.

TABLE 1

Calculated and measured parameter values

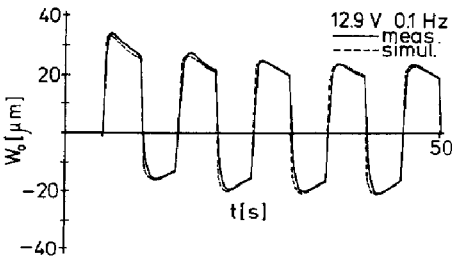
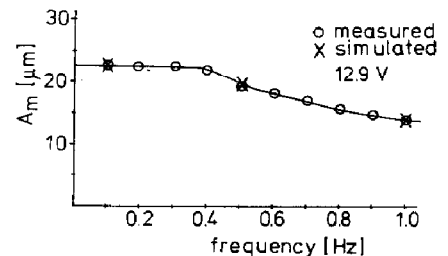
Parameter		Measured (plots)	Calculated (geometry)	Simulated	Units
Volume cavity	V_p		2.7×10^{-8}	2.7×10^{-8}	m^3
Flow conduct. channels	G_{pc}		4.0×10^{-16}		$m^5/N s$
Flow conduct. leak	G_{pl}			8.0×10^{-14}	$m^5/N s$
Relaxation time	τ_p	10		9	s
Relaxation time	τ_p'	11		9	s
Relaxation time	τ_p''	9		9	s
Heat capacity sheet	C_t		4.4×10^{-3}	6.0×10^{-3}	J/K
Heat conductivity, air	G_{ta}		1.2×10^{-2}	1.2×10^{-2}	W/K
Heat conductivity, beams	G_{tb}		1.5×10^{-2}	1.8×10^{-2}	W/K
Response time	τ_t	0.25	0.16	0.20	s
Response time	τ_t'	0.25	0.16	0.20	s

Figure 3 shows w_0 when a block voltage of 12.9 V is applied to the resistor at 0.1 Hz. The system moves towards a steady state with a relaxation time τ_p'' , which is also determined by the volume of the cavity, V_p , and the flow conductance of the leak, G_{pl} ; therefore τ_p'' approximately equals τ_p' and τ_p . After 30 s this steady state is reached with an amplitude A_m of 23 μm . The dashed line corresponds to a simulation.

At higher frequencies the actuator shows a similar behaviour. However, at frequencies exceeding 0.4 Hz, there is no time to reach maximum or minimum w_0 , resulting in a decrease of A_m . Therefore, the maximum of A_m is determined by input power, G_{ta} and G_{tb} , and frequency. In Fig. 4 measured and simulated amplitudes are given as a function of frequency, for a block voltage of 12.9 V.

The temperature rise of the actuator as a whole, as determined by means of the attached thermocouple, appears to be slow and small. For instance, a maximum temperature rise of 10 K is reached after 5 min at an input power of 0.5 W.

In order to determine the influence of the presence of a liquid on the behaviour of the actuator, to check its suitability for application in a pump,

Fig. 3. Deflection w_0 vs. time for a block voltage of 12.9 V and 0.1 Hz.Fig. 4. Measured and simulated amplitudes A_m vs. frequency for a block voltage of 12.9 V.

the membrane was covered with a drop of water. The presence of this water has no measurable effect on w_0 or A_m .

A deflection of 23 μm corresponds to a pressure difference over the membrane of 0.02 atm and a volume stroke of 0.5 μl . Reduction of the thermal conductances G_{ta} and G_{tb} will result in an increase in temperature and pressure build-up.

If the actuator is to operate at frequencies exceeding 0.4 Hz, the relaxation times τ and τ' have to be reduced. This can be achieved by reducing the heat capacity C_t of the suspended thin silicon sheet.

Conclusion

The actuator behaves as expected: results of simulations and measurements coincide within 10%. Hence the simulation model can be used for optimizing the actuator for a required pressure build-up in the cavity and a specific related (free) volume stroke of the membrane.

Important parameters are the response time τ_t (determining maximum frequency f_{max}) and the maximum build-up pressure Δp_{max} . The heat capacity of the suspended thin silicon sheet supporting the aluminium resistor determines, to a great extent, τ_t and f_{max} , while Δp_{max} is mainly determined by input power and thermal conductances. Both Δp_{max} and the induced membrane deflection are sufficient to drive the microminiature pump we have in mind, and can be even further optimized.

The construction of the actuator is simple and it can be fabricated using IC-compatible materials.

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