

**EXACT VALUE OF THE GROUND STATE ENERGY
OF THE LINEAR ANTIFERROMAGNETIC HEISENBERG CHAIN
WITH NEAREST AND NEXT-NEAREST NEIGHBOUR INTERACTIONS**

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It is shown that the ground state energy of the hamiltonian $H = \sum \mathbf{S}_i \cdot \mathbf{S}_{i+1} + \gamma \sum \mathbf{S}_i \cdot \mathbf{S}_{i+2}$ for the linear antiferromagnetic Heisenberg chain with nearest and next-nearest neighbour interactions is equal to $-\frac{3}{2}$ if $\gamma = \frac{1}{2}$.

We consider the linear antiferromagnetic Heisenberg chain with nearest and next-nearest neighbour interactions with hamiltonian

$$H_N(\gamma) = \sum_{i=1}^{N-1} \mathbf{S}_i \cdot \mathbf{S}_{i+1} + \gamma \sum_{i=1}^{N-2} \mathbf{S}_i \cdot \mathbf{S}_{i+2}, \quad (1)$$

where N is the number of spins $\frac{1}{2}$ in the chain. Let the lowest eigenvalue of $H_N(\gamma)$ be $E_N(\gamma)$. We are interested in the energy per spin in the ground state in the limit of large N :

$$E(\gamma) = \lim_{N \rightarrow \infty} E_N(\gamma)/N. \quad (2)$$

$E(0)$ is known exactly [1,2]; it is equal to $1 - 4 \ln 2$ ($= -1.7726, \dots$). For $\gamma \neq 0$ only approximations and upper and lower bounds for $E(\gamma)$ are known [3-8]. Majumdar and Ghosh [3,4] have found that for small finite chains ($N \leq 10$) with an even number of spins and periodic boundary conditions $E_N(\frac{1}{2})$ is equal to $-\frac{3}{2}N$. Moreover they showed that $-\frac{3}{2}N$ is an eigenvalue of $H_N(\frac{1}{2})$ for all even N . The same is true for open chains with N even [5]. Thus $E(\frac{1}{2}) \leq -\frac{3}{2}$. The aim of this note is to show that $E(\frac{1}{2})$ is equal to $-\frac{3}{2}$ and that $E(\gamma)$ takes its maximal value at $\gamma = \frac{1}{2}$.

Divide the chain into $\frac{1}{2}N$ cells of 2 spins. Let the state $|\psi\rangle$ of the chain be the direct product of the states, $2^{-1/2}(|+\rightarrow\rangle - |-\rightarrow\rangle)$ for the cells. Then it is easy to verify that

$$H_N(\frac{1}{2})|\psi\rangle = -\frac{3}{2}N|\psi\rangle, \quad (3)$$

and that

$$\langle\psi|H_N(\gamma)|\psi\rangle = -\frac{3}{2}N. \quad (4)$$

Therefore

$$E(\gamma) \leq -\frac{3}{2}. \quad (5)$$

This shows that $E(\gamma)$ takes its maximal value $-\frac{3}{2}$ at $\gamma = \frac{1}{2}$ if we show that $E(\frac{1}{2}) = -\frac{3}{2}$. To show that $E(\frac{1}{2}) = -\frac{3}{2}$ we calculate a lower bound for $E(\gamma)$.

The hamiltonian H_N can be written as

$$H_N(\gamma) = \sum_{i=1}^{N-2} H_i(\gamma) + \frac{1}{2}\mathbf{S}_1 \cdot \mathbf{S}_2 + \frac{1}{2}\mathbf{S}_{N-1} \cdot \mathbf{S}_N, \quad (6)$$

where

$$H_i(\gamma) = \frac{1}{2}\mathbf{S}_i \cdot \mathbf{S}_{i+1} + \frac{1}{2}\mathbf{S}_{i+1} \cdot \mathbf{S}_{i+2} + \gamma\mathbf{S}_i \cdot \mathbf{S}_{i+2}. \quad (7)$$

Since, in general, the lowest eigenvalue of a sum of operators is not less than the sum of the lowest eigenvalues of these operators, the lowest eigenvalue of $H_N(\gamma)$ is not less than $(N-2)E_i(\gamma) - 3$ if $E_i(\gamma)$ is the lowest eigenvalue of $H_i(\gamma)$. $E_i(\gamma)$ can be calculated immediately; we find

$$\begin{aligned} E_i(\gamma) &= \gamma - 2, & \text{if } \gamma \leq \frac{1}{2}, \\ &= -3\gamma, & \text{if } \gamma \geq \frac{1}{2}. \end{aligned} \quad (8)$$

Since $E_i(\frac{1}{2}) = -\frac{3}{2}$ it follows that $E_N(\frac{1}{2}) \geq -\frac{3}{2}N$. It

follows that $E(\frac{1}{2}) \geq -\frac{3}{2}$ which proves that $E(\frac{1}{2}) = -\frac{3}{2}$.

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