

Mathematical modelling of a flow-injection system with a membrane separation module

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(Received 24th February 1992)

Abstract

A mathematical model for a flow-injection system with a membrane separation module based on the axially dispersed plug flow model was developed. It takes into account the geometrical dimensions and dispersion properties of the main sections of the manifold, the mass transfer in the channels of the separation module and the characteristics of the membrane (thickness and diffusion coefficient within it). The model was solved analytically in the Laplace domain. The inverse transformation was found to give satisfactory results for reactor Peclet numbers less than 120. Otherwise a numerical solution based on the implicit alternating-direction finite difference method was preferred. The adequacy of the model was confirmed experimentally on a flow-injection manifold with a parallel-plate dialysis module. The unknown flow and membrane parameters were determined by curve fitting. The membrane parameters were determined also by steady-state measurements. Fairly good agreement between the dynamic and steady-state results and with results given in the literature was observed, which, together with other experimental results, supported the validity of the model and showed that it can be used successfully for the mathematical description and optimization of flow-injection systems with membrane separation modules. In this connection, the influence of the reactor parameters and the sample volume on the performance of such a system were investigated and conclusions for improving its sensitivity and sample throughput were drawn. Other possible applications of the model are in membrane technology for characterizing of various membranes and in process engineering for investigating the mass transfer in different dialysers.

Keywords Flow-injection, Optimization methods, Dialysis separation, Mathematical modelling, Membrane permeability

Flow-injection manifolds with modules for dialysis or gas diffusion have been shown to be very useful in dealing with samples containing macromolecules, particles or cells by separating them from the analyte prior to detection [1–4]. In this way, malfunctioning of the detector due to processes associated with the species mentioned above (e.g., clogging, light scattering) can be ef-

fectively prevented, thus securing accurate and reproducible results. This makes the flow-injection dialysis systems very advantageous in various fields (e.g., clinical analysis [3,4]) where, in addition to the accuracy and reproducibility of the analysis, a high sample throughput with minimum consumption of reagents is of crucial importance.

The mass-transfer process in the dialysis module of a flow-injection system is determined to a considerable extent by the properties of the membrane (i.e., its thickness and diffusivity). This allows experimental characterization of semi-permeable membranes using such systems [5].

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For a better understanding of the processes taking place in a flow-injection manifold with a dialysis module, an adequate mathematical model is necessary. On the one hand, it will give possibilities for deriving guidelines for optimization of both the design and the operation of these manifolds, thus avoiding the costly and time-consuming “trial and error” approach. On the other hand, it will allow the separation of the contribution of the membrane to the overall mass transfer from that of the flow pattern in both channels of the dialysis module. This is necessary when the flow system is used for membrane characterization or for investigating the mass transfer in the channels.

Theoretical investigations of the mass-transfer process in dialysers based on fundamental physical laws (e.g., the Navier–Stokes equations and Fick’s second law) have been limited mainly to the case of fully developed laminar flow in both channels and steady-state mass transfer. The reason for this situation is that industrial mass exchangers and haemodialysers used as artificial kidneys work under such conditions and the overall effect of the transient periods (e.g., start-up, shut-down, power surge and pump failure) is usually negligible. However, in flow-injection manifolds with dialysis separation, the transient mass transfer becomes the process governing the performance. The exact mathematical description must include the time dependence of the mass-transfer process, which greatly complicates the corresponding mathematical solution.

Bernhardsson et al [6] used the approach mentioned above for the modelling of a flow-injection manifold assumed to consist of a dialysis module only and to operate under steady-state conditions. The latter assumption does not permit the calculation of the concentration profiles at the outlets of both channels under flow-injection conditions because they are the result of transient mass transfer. The model was successfully used for determining the membrane permeability [6,7].

The overall performance of a flow-injection manifold with a dialysis module depends on the processes taking place in all the separate sections (e.g., injection valve, reactor). Because of the

complexity of the real hydrodynamic regime in the different sections, the development of a general mathematical model based on fundamental physical laws is exceedingly difficult and in most instances virtually impossible.

A possible approach for predicting the outlet concentration profiles of a given dialysis module is to find its impulse–response function under certain experimental conditions. This can be performed by deconvolution of the concentration profiles obtained under identical experimental conditions in a flow-injection manifold with and without the dialysis module [8,9]. However, it is difficult to relate the impulse–response function obtained in this way to the parameters that most affect the behaviour of the dialyser and to extend the results to other sets of parameters values.

The approach used in chemical engineering for the mathematical description of such complex flow systems is based on the application of the so-called hydraulic models (e.g., tanks-in-series, dispersion, combined and empirical models) [10,11]. In the modelling of flow-injection manifolds, the tanks-in-series and the dispersion models have been most frequently used [1,2]. For the mathematical description of flow-injection systems with membrane separation, the tanks-in-series model without back-mixing and with equal size of the tanks was used [12] and gave satisfactory qualitative agreement with the experimental results. However, most of the flow-through sections in a flow-injection manifold are tubular and for that reason the dispersion models appear to be closer to the real physical situation. Among them, the axially dispersed plug flow model [10] deserves special attention because it is simpler from a mathematical point of view and it has been already successfully utilized in the modelling of single-line flow-injection systems [13–17]. The flow pattern prevailing in these systems is laminar and another advantage offered by the axially dispersed plug flow model is the possibility of calculating the laminar Peclet number in both their tubular [18,19] and parallel-plate [20] sections. For parallel-plate dialysis modules, a theoretical relationship for determining the laminar flow mass-transfer coefficients in both channels has also been derived [21].

TABLE 1

Symbols and definitions ^a

<i>a</i>	Half of a channel height or radius of a tubular section (m)
<i>A</i>	Integrational constants (Appendix A)
<i>B</i>	Integrational constants (Appendix A)
<i>c</i>	Concentration (mol m ⁻³)
<i>c</i> ₀	Initial tracer concentration (mol m ⁻³)
<i>c</i> _{av}	= $c_0 V_{inj} / V_t$ Average concentration (mol m ⁻³)
<i>C</i>	= c / c_{av} (under flow-injection conditions) or = c / c_0 (under steady-conditions) Dimensionless concentration
\bar{C}	Laplace transform of <i>C</i>
<i>D</i>	Molecular diffusion coefficient (m ² s ⁻¹)
<i>D</i> _L	Axial dispersion coefficient (m ² s ⁻¹)
<i>e</i>	Peak broadening (s)
<i>E</i>	= $e v_D / V_t$ Dimensionless peak broadening
<i>f</i>	Coefficient defined in Appendix A
<i>F</i>	= c_m / c_{Dm} or c_m / c_{Am} Distribution coefficients
<i>g</i>	Coefficient defined in Appendix A
<i>h</i>	Coefficient defined in Appendix A
<i>H</i>	= ka_D / D_m Dimensionless group
<i>k</i>	Mass-transfer coefficient (m s ⁻¹)
<i>K</i>	= $k V_t / 2 a v_D$ Dimensionless mass-transfer coefficient
<i>K</i>	Number of θ subintervals (Appendix B)
<i>L</i>	= x_2 Characteristic length (m)
<i>M</i>	Coefficients defined in Appendix A
<i>M</i>	Number of <i>Y</i> subintervals (Appendix B)
<i>N</i>	Number of <i>X</i> subintervals (Appendix B)
<i>p</i>	Laplace complex variable
<i>P</i>	= $\bar{u} L / D_L$ Peclet number
<i>P</i> _m	= $a_D^2 v_D / D_m V_t$ Dimensionless group
<i>q</i>	Coefficients defined in Appendix A
<i>Q</i>	Coefficient defined in Appendix A
<i>r</i>	Integrational constants (Appendix A)
<i>s</i>	Cross-sectional area of a flow section (m ²)
<i>S</i>	Coefficient defined in Appendix A
<i>t</i>	Time (s)
\bar{u}	Mean linear flow-rate (m s ⁻¹)
<i>v</i>	Volumetric flow-rate (m ³ s ⁻¹)
<i>V</i>	= sL Volume (m ³)
<i>V</i> _{inj}	Injected volume of tracer (m ³)
<i>V</i> _t	= $\pi a_D^2 x_1 + 2 a_D w (x_2 - x_1)$ Characteristic volume
<i>w</i>	Width of the channels of the membrane separation module (m)
<i>x</i>	Axial distance (m)
<i>X</i>	= x / L Dimensionless axial distance
<i>y</i>	Transverse distance (m)
<i>Y</i>	= y / a_D Dimensionless transverse distance
<i>z</i>	Roots of Eqn. A20

Greek letters

α	Coefficient defined in Appendix A
β	The only negative root of Eqn. C3
γ	= V_t / V Coefficients
δ	Half-width of the membrane (m)
Δ	= δ / a_D Dimensionless half-width of the membrane
ΔX	<i>X</i> increment (Appendix B)
ΔY	<i>Y</i> increment (Appendix B)

TABLE 1 (continued)

Greek letters

$\Delta\theta$	θ increment (Appendix B)
θ	$= tv_D/V_1$ Dimensionless time
μ	Statistical moment of the tracer response curve about the origin
σ^2	Variance of the tracer response curve
τ	$= Dx/(\bar{u}a^2)$ Fourier number
ϕ	Integrational variable
Ψ	Coefficients defined by Eqn C4
ω	Coefficients defined by Eqn C5

^a Subscripts f, i, r and a refer to the fore-section, the injection section, the reactor and the after-section, respectively, D, A and m refer to the donor channel, the acceptor channel and the membrane, respectively, Dm and Am refer to the donor stream/membrane and the acceptor stream/membrane interfaces, respectively, max refers to maximum concentration. Superscript * refers to the acceptor line.

The aim of this paper is to outline a mathematical model of a flow-injection system with a membrane separation module based on the axially dispersed plug flow model, its experimental verification and the formulation of guidelines for improving its performance.

Among the various types of dialysers used in industry and medical practice and for analytical purposes (i.e., parallel-plate, tubular and spiral-plate dialysers), parallel-plate dialysers have been used most frequently in flow-injection manifolds. For this reason, the experimental verification of the model was performed on a flow-injection manifold with a parallel-plate dialysis module.

DEVELOPMENT AND SOLUTION OF THE MATHEMATICAL MODEL

Description of the model

A flow-injection manifold with a membrane separation module which may have arbitrary geometry can be presented schematically as consisting of one donor and one acceptor line (Fig 1). The donor line includes the following flow-through sections connected in series: the fore-section, which connects the injection device with the reservoir of the carrier solution, the injection section (usually an injection valve), the reactor, which could be a straight, coiled, packed-bed, single-bead string or knitted tube, the donor channel inside the membrane separation module, and the after-section, which is the section of the

flow system leading to waste. Similarly, the acceptor line consists of fore-section, acceptor channel and after-section. The measuring cell is assumed to be downstream of the acceptor channel. Its volume is usually negligible in comparison with the volume of the reactor and the two channels inside the membrane separation module. For this reason, it is not included as a separate element in the schematic representation of the acceptor line. The additional assumptions on which the model is built are the following: (i) the dispersion of the analyte in all sections of the manifold can be described by the axially dispersed plug flow model (Eqn 1), (ii) the fore- and after-sections are tubular and infinitely long, so that the so-called "end effects" [22] can be neglected and Taylor's theory [18,19] can be used for calculating their Peclet

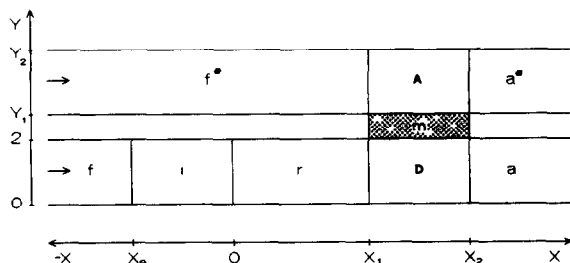


Fig 1 Scheme of flow-injection manifold with membrane separation module. f, i, r and a refer to the fore-section, the injection section, the reactor and the after-section, respectively, D and A refer to the donor and the acceptor channels, respectively, and * refers to the acceptor line, $Y_1 = 2(1 + \Delta)$, $Y_2 = 2(1 + \Delta + a_A/a_D)$

numbers (Eqn 2), and the mass transport within the membrane is caused by transverse Fickian diffusion only (All symbols are defined in Table 1)

$$\frac{\partial C}{\partial \theta} = \frac{1}{P} \frac{\partial^2 C}{\partial X^2} - \frac{\partial C}{\partial X} \quad (1)$$

$$P = 48DL / (\bar{u}a^2) \quad (2)$$

The flow system shown in Fig 1 can be described by the following set of partial differential equations in dimensionless quantities and variables similar to Eqn 1

For the donor line

$$\frac{\partial C_f}{\partial \theta} + \gamma_f \frac{\partial C_f}{\partial X} - \frac{\gamma_f}{P_f} \frac{\partial^2 C_f}{\partial X^2} = 0 \quad X \leq X_e \quad (3)$$

$$\frac{\partial C_1}{\partial \theta} + \gamma_1 \frac{\partial C_1}{\partial X} - \frac{\gamma_1}{P_1} \frac{\partial^2 C_1}{\partial X^2} = 0 \quad X_e \leq X \leq 0 \quad (4)$$

$$\frac{\partial C_r}{\partial \theta} + \gamma_r \frac{\partial C_r}{\partial X} - \frac{\gamma_r}{P_r} \frac{\partial^2 C_r}{\partial X^2} = 0 \quad 0 \leq X \leq X_1 \quad (5)$$

$$\frac{\partial C_D}{\partial \theta} + \gamma_D \frac{\partial C_D}{\partial X} - \frac{\gamma_D}{P_D} \frac{\partial^2 C_D}{\partial X^2} + K_D(C_D - C_{Dm}) = 0 \quad X_1 \leq X \leq X_2 \quad (6)$$

$$\frac{\partial C_a}{\partial \theta} + \gamma_a \frac{\partial C_a}{\partial X} - \frac{\gamma_a}{P_a} \frac{\partial^2 C_a}{\partial X^2} = 0 \quad X \geq X_2 \quad (7)$$

For the membrane

$$\frac{\partial C_m}{\partial \theta} - \frac{1}{P_m} \frac{\partial^2 C_m}{\partial Y^2} = 0 \quad 2 \leq Y \leq 2 + 2\Delta \quad (8)$$

For the acceptor line

$$\frac{\partial C_f^*}{\partial \theta} + \gamma_f^* \frac{\partial C_f^*}{\partial X} - \frac{\gamma_f^*}{P_f^*} \frac{\partial^2 C_f^*}{\partial X^2} = 0 \quad X \leq X_1 \quad (9)$$

$$\frac{\partial C_A}{\partial \theta} + \gamma_A \frac{\partial C_A}{\partial X} - \frac{\gamma_A}{P_A} \frac{\partial^2 C_A}{\partial X^2} - K_A(C_{Am} - C_A) = 0 \quad X_1 \leq X \leq X_2 \quad (10)$$

$$\frac{\partial C_a^*}{\partial \theta} + \gamma_a^* \frac{\partial C_a^*}{\partial X} - \frac{\gamma_a^*}{P_a^*} \frac{\partial^2 C_a^*}{\partial X^2} = 0 \quad X \geq X_2 \quad (11)$$

The coefficients γ make it possible to use the equations given above for the description of flow systems consisting of sections with various cross-sections. When all cross-sections have equal diameters, the coefficients γ are unity

The initial and boundary conditions are the following

For the donor line

$$C_f(X, 0) = C_r(X, 0) = C_D(X, 0) = C_a(X, 0) = 0, \quad (12a)$$

$$C_1(X, 0) = V_i/V_{im} \quad (12a)$$

$$C_f(X_e^-, \theta) = C_1(X_e^+, \theta) \quad (12b)$$

$$C_f(X_e^-, \theta) - \frac{1}{P_f} \frac{\partial C_f(X_e^-, \theta)}{\partial X} = C_1(X_e^+, \theta) - \frac{1}{P_1} \frac{\partial C_1(X_e^+, \theta)}{\partial X} \quad (12c)$$

$$C_1(0^-, \theta) = C_r(0^+, \theta) \quad (12d)$$

$$C_1(0^-, \theta) - \frac{1}{P_1} \frac{\partial C_1(0^-, \theta)}{\partial X} = C_r(0^+, \theta) - \frac{1}{P_r} \frac{\partial C_r(0^+, \theta)}{\partial X} \quad (12e)$$

$$C_r(X_1^-, \theta) = C_D(X_1^+, \theta) \quad (12f)$$

$$C_r(X_1^-, \theta) - \frac{1}{P_r} \frac{\partial C_r(X_1^-, \theta)}{\partial X} = C_D(X_1^+, \theta) - \frac{1}{P_D} \frac{\partial C_D(X_1^+, \theta)}{\partial X} \quad (12g)$$

$$C_D(X_2^-, \theta) = C_a(X_2^+, \theta) \quad (12h)$$

$$C_D(X_2^-, \theta) - \frac{1}{P_D} \frac{\partial C_D(X_2^-, \theta)}{\partial X} = C_a(X_2^+, \theta) - \frac{1}{P_a} \frac{\partial C_a(X_2^+, \theta)}{\partial X} \quad (12i)$$

$$C_f(-\infty, \theta) = \text{finite} \quad (12j)$$

$$C_a(\infty, \theta) = \text{finite} \quad (12k)$$

For the membrane

$$C_m(Y, 0) = 0 \quad (13a)$$

$$\frac{\partial C_m(2, \theta)}{\partial Y} = -H_D [C_D(X, \theta) - C_{Dm}(X, \theta)] \quad (13b)$$

$$\frac{\partial C_m(2 + 2\Delta, \theta)}{\partial Y} = H_A [C_{Am}(X, \theta) - C_A(X, \theta)] \quad (13c)$$

For the acceptor line

$$C_f^*(X, 0) = C_A(X, 0) = C_a^*(X, 0) = 0 \quad (14a)$$

$$C_f^*(X_1^-, \theta) = C_A(X_1^+, \theta) \quad (14b)$$

$$C_f^*(X_1^-, \theta) - \frac{1}{P_f^*} \frac{\partial C_f^*(X_1^-, \theta)}{\partial X} = C_A(X_1^+, \theta) - \frac{1}{P_A} \frac{\partial C_A(X_1^+, \theta)}{\partial X} \quad (14c)$$

$$C_A(X_2^-, \theta) = C_a^*(X_2^+, \theta) \quad (14d)$$

$$C_A(X_2^-, \theta) - \frac{1}{P_A} \frac{\partial C_A(X_2^-, \theta)}{\partial X} = C_a^*(X_2^+, \theta) - \frac{1}{P_a^*} \frac{\partial C_a^*(X_2^+, \theta)}{\partial X} \quad (14e)$$

$$C_f^*(-\infty, \theta) = \text{finite} \quad (14f)$$

$$C_a^*(\infty, \theta) = \text{finite} \quad (14g)$$

The superscripts $-$ and $+$ refer to the upstream and downstream side, respectively, of a given interface between two adjacent flow sections

The empirical parameters of the model are the Peclet numbers of the different sections and the mass-transfer coefficients in both channels of the membrane separation module. The other parameters of the model are either fundamental constants (e.g., diffusion coefficient in the membrane) or physical parameters, some of which can be measured (e.g., geometrical dimensions, flow-rates). It should be pointed out that the membrane when contacting with solutions usually swells and its thickness under working conditions

cannot be measured easily. There are theoretical and empirical relationships for calculating the laminar Peclet numbers of certain type of reactors (e.g., straight and coiled tubes), valid usually for the case of diffusion-controlled dispersion [16]. In most of the flow-injection systems an injection valve is used. In this instance the injection section includes the sample loop and the internal bores of the valve, thus having a rather complex geometry. Therefore, its Peclet number can be determined experimentally only. However, it has been found that if the volume of the injection section is considerably smaller than the reactor volume (e.g., less than 10% of it), which is frequently the case, the Peclet number of the injection section can be assumed to be equal to that of the reactor [17]. For volumes of the injection section of less than 3% a delta-function approximation can be used, which considerably simplifies the solution of the model [17]. Theoretical relationships have been derived for calculating the Peclet number and the mass-transfer coefficients in both channels of parallel-plate dialysers under laminar flow conditions [20,21].

Solution of the model

Equations 3-11 were solved by means of the Laplace transform technique

$$\bar{C}(X, Y) = \int_0^\infty C(\theta, X) \exp[-p\theta] d\theta \quad (15)$$

In the Laplace domain these equations reduce to ordinary linear differential equations of the second order, which were solved analytically (Appendix A). Interesting points for monitoring the concentration are the inlet and outlet of the donor channel and the outlet of the acceptor channel in the case of both semi-permeable and impermeable membranes. The corresponding Laplace domain solutions of mass transfer in the membrane separation module are presented in Appendix A for three different cases: (i) overall transfer is governed by transfer in the channels and in the membrane; (ii) overall transfer is governed by transfer in the membrane only (i.e., infinitely high mass-transfer coefficients), and (iii) only transfer in the channels is of importance

(i.e., infinitely thin membrane or infinitely high diffusion coefficient within it) The time domain solutions were obtained by numerical inverse transformation of the corresponding Laplace domain solutions by expansion of the latter into Fourier sine series and subsequent analytical inverse transformation [23]

For high Peclet numbers of the reactor and the channels of the separation module (e.g., $P > 120$), the numerical technique described above does not give satisfactory results. The concentration curves are lower and exhibit oscillating fronts and tails. Taking into consideration that at high Peclet numbers the "end effects" are negligible [22], it can be assumed that the flow pattern in the membrane separation module does not affect the overall dispersion process in the reactor. This allows one to describe the concentration profile at the inlet of the donor channel as if the reactor is infinitely long. For a straight tube reactor in the case of diffusion-controlled dispersion (i.e., $\tau > 0.7$, Eqn 2) and small sample volume, so that the injection section can be considered as part of the reactor, the corresponding analytical solution of Eqn 5 will be the following [24]

$$C(\theta, X_1) = \frac{1}{2} \left\{ \operatorname{erf} \left[\frac{\gamma_r \theta - X_1}{2(\gamma_r \theta / P_r)^{1/2}} \right] + \operatorname{erf} \left[\frac{X_1 - X_e \gamma_r / \gamma_1 - \gamma_r \theta}{2(\gamma_r \theta / P_r)^{1/2}} \right] \right\} \quad (16)$$

where erf is the error function

Equation 16 defines the input concentration profile for the membrane separation module. The equations describing the processes taking place in the two channels of the membrane separation module and in the membrane itself (i.e., Eqns 6, 8 and 10) can be solved by the implicit alternating-direction finite-difference method [25]. The corresponding finite-difference equations are given in Appendix B. The concentration profiles monitored downstream of the outlet of the donor (Eqn 17) and the acceptor (Eqn 18) channels (Fig 1) can be calculated on the basis of the

convolution theorem

$$C(\theta, X) = \left(\frac{P_a}{4\pi\gamma_a} \right)^{1/2} \int_0^\theta C_D(\theta, X_2)(\theta - \phi)^{-1/2} \times \exp \left[-P_a \frac{(X - X_2 - \gamma_a(\theta - \phi))^2}{4\gamma_a(\theta - \phi)} \right] d\phi \quad (17)$$

$$C(\theta, X) = \left(\frac{P_a^*}{4\pi\gamma_a^*} \right)^{1/2} \int_0^\theta C_A(\theta, X_2)(\theta - \phi)^{-1/2} \times \exp \left[-P_a^* \frac{(X - X_2 - \gamma_a^*(\theta - \phi))^2}{4\gamma_a^*(\theta - \phi)} \right] d\phi \quad (18)$$

where $C_D(\theta, X_2)$ and $C_A(\theta, X_2)$ are the numerical solutions of Eqns 6, 8 and 10, and Eqn 19 is the solution of the axially dispersed plug flow model in the case of an infinitely long tube and delta-function injection at $X = X_2$

$$C(\theta, X) = \left(\frac{P}{4\pi\gamma\theta} \right)^{1/2} \times \exp \left[-P \frac{(X - X_2 - \gamma\theta)^2}{4\gamma\theta} \right] \quad (19)$$

EXPERIMENTAL

Flow-injection manifold

The main geometrical dimensions of the experimental flow system used in this work, and presented schematically in Fig 2, are summarized in Table 2

A constant flow-rate in both the donor and the acceptor lines was maintained by a low-pulsation computer-controlled peristaltic pump (Minipuls, Gilson). To reduce further the undesirable effects caused by pulsation and to remove the air bubbles from the donor and the acceptor streams, pulse dampers were installed downstream of the

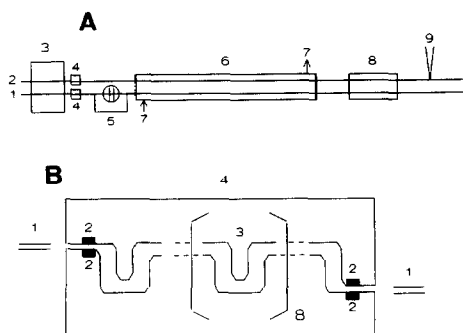


Fig 2 (A) Scheme of the experimental manifold 1 = Donor line, 2 = acceptor line, 3 = peristaltic pump, 4 = pulse dampers, 5 = injection valve, 6 = thermostating Perspex tube, 7 = thermostating water, 8 = dialysis module, 9 = thermocouple (B) Scheme of one of the halves of the dialysis module 1 = PTFE tubes, 2 = platinum electrodes, 3 = element of the channel, repeated eight consecutive times, 4 = Perspex body

pump (Fig 2) One of their walls consists of a gas-permeable membrane made by stretching a PTFE foil The flow-rates in both lines were determined for each experiment by collecting and weighing their effluents over intervals of 5 min The total lengths of the tubings of both lines were chosen to be almost equal to prevent a pressure difference between the two channels in the dialysers, which may cause deformation of the membrane, making the mass-transfer process more complicated The flow-rates in the two lines of the manifold during all the experiments did not differ by more than 1%

An automatic computer-controlled rotary injection valve (Universal, Anachem) with an external sample loop was used The latter was a bent piece of PTFE tubing It was assumed that the whole sample volume filling the external loop and some of the internal bores of the valve were situated in the external loop only On this basis, the length of the external loop, which corresponds in this instance to the injection section in the model (Fig 1), was recalculated (Table 2)

The fore-sections, the after-sections and the reactor were PTFE tubes Their inside diameters were determined by the volume of doubly distilled water with which they can be filled The reactor was a straight tube The sections of the

fore- and after-sections adjacent to the valve and the dialysis module were also straight tubes Their lengths were chosen in such a way that the so-called "end effects" could be neglected [22]

A parallel-plate dialyser made of two Perspex halves with a meander-type rectangular groove (Fig 2A) was used The membrane was placed between the two halves before they were screwed together The bottom of each groove was roughened by milling into small circles, thus creating turbulence promoters for intensifying the mass transfer in the channels Short rectangular channels ($5 \mu\text{l}$ each) connected the grooves of the dialyser with the PTFE tubes by ferrules and screws Their volumes were considered in the calculations as parts of the corresponding adjacent flow sections (Fig 1) Platinum plates ($0.8 \times 0.8 \text{ mm}$) were installed parallel to each other in the bores next to the grooves (Fig 2a) They allowed the conductivity (PW9509 digital conductivity meter, Philips) of the influent and effluent of both the donor and the acceptor channels of the dialyser to be monitored

Two types of hydrophilic membrane were used in the experiments, i.e., a PTFE membrane impermeable to potassium chloride and a permeable cellulose hydrate Cuprophane (Enka Glanzstoff) membrane

The flow system was kept at $20.0 \pm 0.1^\circ\text{C}$ The thermostating system consisted of a thermostat (P M Tamson) and a straight Perspex tube in which the reactor and the fore-section of the acceptor line were installed and through which water at the above temperature was constantly

TABLE 2

Geometrical dimensions of the manifold

Section	Donor line		Acceptor line	
	Length (m)	Diameter (mm)	Length (m)	Diameter (mm)
Fore-section	0.500	0.5304	2.950	0.5304
Injection section	0.115	0.5304		
Reactor	2.487	0.5304		
Channel ^a	0.166	0.2405 ^b	0.166	0.2590 ^b
After-section	0.540	0.5304	0.450	0.5304

^a The width of both channels is 0.0030 m ^b Height of the channels

circulating. The source bottles containing the carrier solutions for the donor and the acceptor lines were placed in the basin of the thermostat. The temperature of the acceptor stream was controlled by a platinum-platinum/rhodium thermocouple connected to a microvoltmeter (3478A Multimeter, Hewlett Packard) (Fig. 2).

The experimental manifold described above was operated by an Apple IIG computer. The software was written in Fysforth version 0.3. The data readings from the conductivity meter and the microvoltmeter were collected and transferred to an IBM-compatible PC for further processing.

Reagents

Standard solutions of potassium chloride in doubly distilled water with concentrations in the range 1.0×10^{-4} – 1.0×10^{-2} M were used for injection, calibration and as carrier solutions in the steady-state experiments.

Procedure

The stimulus-response technique [10] was used for identification of the unknown parameters in the model outlined in this paper under flow-injection conditions. The tracer injected by the valve was a 1.0×10^{-2} M solution of potassium chloride. The carrier solutions in both lines, if not stated otherwise, were doubly distilled water.

Three series of experiments were performed as follows. In the first series, the membrane (PTFE foil) installed in the dialysis module was impermeable to the tracer (potassium chloride). The response curves at the inlet and outlet of the donor channel were monitored at four different flow-rates in the range 0.1–1.0 ml min⁻¹. For investigating the acceptor channel its position was interchanged with that of the donor channel, thus connecting it to the reactor. The same experimental procedure was performed again.

In the second series of experiments, the impermeable PTFE foil was replaced by Cuprophane membrane, permeable to the tracer. The response curves at the inlet and outlet of the donor channel and at the outlet of the acceptor channel were recorded at different flow-rates in the range 0.1–0.5 ml min⁻¹. For higher flow-rates the mass

transfer under the working conditions used became negligible.

In the third series, doubly distilled water was maintained as the carrier solution in the acceptor line while the carrier solution in the donor line was 1.6×10^{-3} M potassium chloride solution. The steady-state concentrations at the outlets of both channels of the dialyser and at the inlet of the donor channel were measured at different flow-rates ranging from 0.1 to 0.5 ml min⁻¹ as in the second series of experiments.

Calibration graphs for all four conductivity detectors were obtained on the basis of fourteen standard solutions of potassium chloride (0.0–2.2 $\times 10^{-3}$ M). The relationship between concentration and conductance was described by polynomial regression equations of the second power. It was found that the flow-rate had no effect on the calibration graphs.

Processing of the experimental response curves

The unknown parameters of the model are the Peclet numbers in all eight flow sections (Fig. 1), the mass-transfer coefficients of the donor and the acceptor channels, the diffusion coefficient and the thickness of the Cuprophane membrane. As already pointed out, the lengths of the fore- and after-sections were chosen such that their Peclet numbers (Eqn. 2) can be determined according to Taylor's theory [18,19]. The volume of the tracer injected in each experiment was less than 5% of the characteristic volume (V_c) (Table 2). In such a case the overall dispersion process should depend only slightly on the Peclet number of the injection section [17] and the latter can be assumed to be equal to that of the reactor. The flow-rates used in the experiments corresponded to Fourier numbers of the reactor (τ_r) well above 0.7, which made it possible to apply Eqn. 2 for the calculation of the Peclet number. On the basis of the above considerations, the eight unknown Peclet numbers could be reduced to three unknown parameters, i.e., the Peclet numbers of the donor and the acceptor channels and the diffusion coefficient of potassium chloride in doubly distilled water at 20°C. The last parameter was necessary for calculating the remainder of the Peclet numbers of the flow system according

to Eqn 2. In fact, these three parameters are the only unknown parameters of the model under the assumptions made above in the case of an impermeable membrane (i.e., first series of experiments). From the response curves detected in the donor line before and after the donor or the acceptor channels, it was possible to determine the diffusion coefficient of the tracer in the liquid phase and the axial dispersion coefficients of both channels at different flow-rates. The other unknown parameters of the model, i.e., the mass transfer in the channels and within the membrane, were determined from the response curves monitored at the outlet of the donor and the acceptor channels in the second series of experiments (with the Cuprophane membrane). The results obtained were used to derive empirical relationships for the calculation of the mass-transfer and axial dispersion coefficients for both channels.

The values of the steady-state concentrations at the outlets of both channels measured at different flow-rates can be used for calculating the permeability of the Cuprophane membrane. The appropriate equations are given in Appendix C. For each flow-rate, the permeability of the membrane was determined as the average of the permeabilities based on the steady-state concentrations measured in donor and acceptor channels (Appendix C). In these calculations, the mass-transfer and axial dispersion coefficients (Peclet numbers) of the two channels of the dialyser were calculated by the relationships derived earlier. The proper functioning of the manifold was controlled by monitoring the tracer concentration at the inlet of the donor channel, which should be constant during all steady-state experiments.

Curve fitting was utilized to determine the unknown parameters when the tracer response curves were processed (i.e., in the first two series of experiments). A simplex optimization method based on the algorithm of Nelder and Mead [26] was used. The function minimized by this procedure was the square root of the mean squared error between experimental and theoretical response curves. Only those parts of the response curves where the concentration was greater than 5% of the maximum concentration were used in

the calculations, thus saving computational time and excluding the less informative sections in the front and the tail of the experimental response curves. The zeroth and first statistical moments about the origin (Eqn 20, $j = 0$ and $j = 1$) and the second moment about the mean (Eqn 21) of the experimental response curves and of their best theoretical fits were calculated and compared.

$$\mu_j = \int_0^\infty C \theta^j d\theta \quad (20)$$

$$\sigma^2 = \int_0^\infty C (\mu_1 - \theta)^2 d\theta = \mu_2 - \mu_1^2 \quad (21)$$

It is possible to calculate the statistical moments directly from the Laplace domain solution of the model (Appendix A) without the necessity of performing its inverse transformation by the following relationship, proposed by Van der Laan [27]

$$\mu_j = (-1)^{-j} (d^j \bar{C} / d p^j)_{p \rightarrow 0} \quad (22)$$

The zeroth moments for the donor and the acceptor channels define the integral amount of tracer monitored in them. They are interrelated by the following equation

$$\mu_{0D} + (v_A/v_D) \mu_{0A} = 1 \quad (23)$$

This relationship was used for controlling the manifold for malfunctioning during the experiments. The first moment about the origin (Eqn 20, $j = 1$), known as the mean, defines the centre of gravity of the tracer response curve. The second moment about the mean, known as the variance, (Eqn 21) characterizes the width of the tracer response curve and can be used for determining the sample throughput of the manifold.

Computer programs

The numerical procedures outlined above for processing of the tracer response curve and for solving the model were programmed in QuickC (Version 2, Microsoft) and the corresponding programs were run on an IBM-compatible PC.

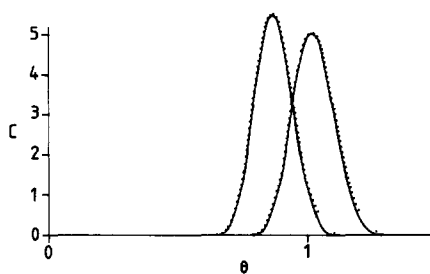


Fig 3 Comparison of experimental tracer response curves (solid lines) measured at $0.133 \text{ ml min}^{-1}$ before and after the donor channel and their best theoretical fits (dotted lines)

RESULTS AND DISCUSSION

Determination of the diffusion coefficient of the tracer in the flow system and its axial dispersion coefficients in the channels of the dialyser

Figure 3 illustrates the agreement between the experimental response curves from the first series of experiments (with the impermeable membrane) and their best theoretical fits

The diffusion coefficient of potassium chloride obtained from the Peclet number of the reactor (Eqn 2) by the curve-fitting procedure, outlined earlier was found to be $1.72 \times 10^{-9} \text{ m}^2 \text{ s}^{-1}$, which was fairly close to earlier results [15] and to literature data ($1.77 \times 10^{-9} \text{ m}^2 \text{ s}^{-1}$ at infinite dilution [28]). Very slight variations in this value (less than 3%) were observed for different flow-rates in the range $0.1-1.0 \text{ ml min}^{-1}$. Because the geometry of the dialyser channels is complicated, no theoretical relationships similar to Eqn 2 for

calculating the corresponding axial dispersion coefficients or Peclet numbers exist. It was found that the flow-rate dependence of the axial dispersion coefficients can be described successfully (Fig 4) by the following empirical equations

$$D_{L_D} = 1.531 \times 10^6 v_D^{1.2692} \quad (24)$$

$$D_{L_A} = 9.865 \times 10^4 v_A^{1.1414} \quad (25)$$

Determination of the mass-transfer coefficients in the channels of the dialyser, the thickness of the membrane and the diffusion coefficient of the tracer within it

The tracer response curves obtained in the second series of experiments were processed by curve-fitting for determining the values of the following parameters of the model K_D , K_A , P_m and Δ . From the last two parameters the thickness of the membrane and the effective diffusion coefficient of potassium chloride within it at 20°C can be easily calculated. In the calculations all the Peclet numbers in the model were calculated by using Eqn 2 and Eqns 24 and 25. In the second series of experiments, the flow-rate in both lines of the manifold was varied ($0.1-0.5 \text{ ml min}^{-1}$), thus changing the hydrodynamic conditions in both channels of the dialysis module. It can be expected that the corresponding mass-transfer coefficients will vary also with the flow-rate while the thickness of the membrane (2δ) and the diffusion coefficient of the tracer within it (D_m) should remain unchanged. The experimental results showed that 2δ and D_m obtained for four different flow-rates were within 3% of their mean values (Table 3), i.e., $2.808 \times 10^{-5} \text{ m}$ and $7.915 \times 10^{-11} \text{ m}^2 \text{ s}^{-1}$, respectively. This result confirms the validity of the model. The

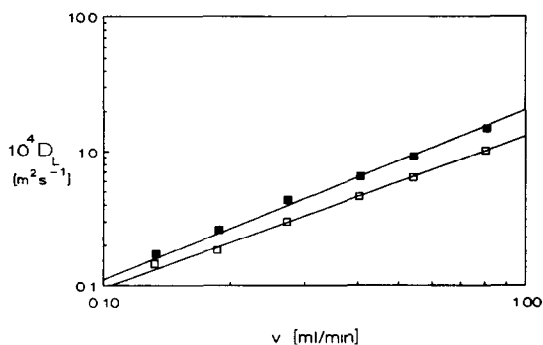


Fig 4 Axial dispersion coefficient versus flow-rate in ml min^{-1} for the (■) donor and (□) acceptor channels

TABLE 3

Curve-fitting results for 2δ and D_m

v_D (ml min^{-1})	2δ (μm)	$10^{11} D_m$ ($\text{m}^2 \text{ s}^{-1}$)
0.1016	27.62	7.747
0.1917	28.08	8.141
0.2803	28.30	7.892
0.4153	28.32	7.880
Average	28.08	7.915

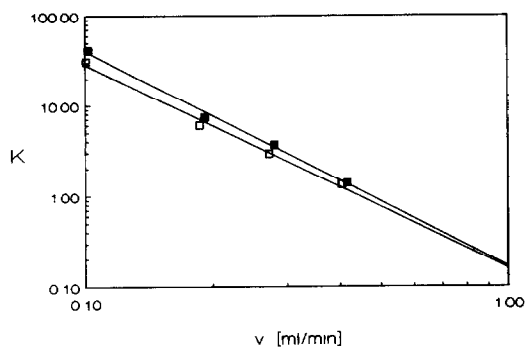


Fig 5 Mass-transfer coefficient versus flow-rate in ml min^{-1} for the (■) donor and (□) acceptor channels

thickness of the membrane (2δ) before contacting with water (normal state) is stated by the producer to be 1.3×10^{-5} m. Its increase during the transition from the normal to the wet state due to swelling is reported to amount to a factor of 1.9. This will result in a 2δ value of 2.48×10^{-5} m, which is close to the value obtained in the present investigation. This fact could be considered as further proof for the adequacy of the model outlined earlier. The values obtained for the mass-transfer coefficients of the channels for different flow-rates were found to obey fairly well the following relationships (Fig 5)

$$K_D = 4.854 \times 10^{-20} v_D^{-2.382} \quad (26)$$

$$K_A = 3.465 \times 10^{-19} v_A^{-2.270} \quad (27)$$

The satisfactory agreement between the experimental response curves and their best theoretical fits is illustrated in Fig 6

Determination of the permeability of the membrane under steady-state conditions

The steady-state concentrations measured at the outlets of the donor and the acceptor channels for different flow-rates are presented in Fig 7

Under steady-state conditions, the thickness of the membrane and the diffusion coefficient of the tracer within it cannot be determined separately from each other, because only the ratio of these two parameters appears in the equations describing the steady-state mass transfer through the membrane (Appendix C). For this reason, the

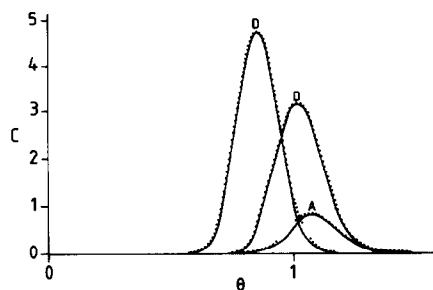


Fig 6 Experimental tracer response curves (solid lines) measured at $v_D = 0.190 \text{ ml min}^{-1}$ and $v_A = 0.187 \text{ ml min}^{-1}$ before the donor channel and after the donor and the acceptor channels and their best model fits (dotted lines). D and A refer to the donor and acceptor lines, respectively

experimental results obtained in this series of experiments were used for determining the ratio $D_m/2\delta$, which is often referred to as the permeability of the membrane. The value of the permeability, averaged over the whole range of flow-rates, was found to be $2.35 \times 10^{-6} \text{ m s}^{-1}$, which does not differ substantially from the corresponding value obtained under flow-injection conditions, i.e., $2.82 \times 10^{-6} \text{ m s}^{-1}$. This result is further proof of the validity of the model. Permeability data for Cuprophane membranes with thickness 1.15×10^{-6} and 1.35×10^{-6} m, when a 0.01 M solution of sodium chloride at 20°C was used, were provided by the manufacturer. According to the manufacturer, the membrane used in the present experiment had a thickness of 1.27×10^{-6} m. For this membrane the perme-

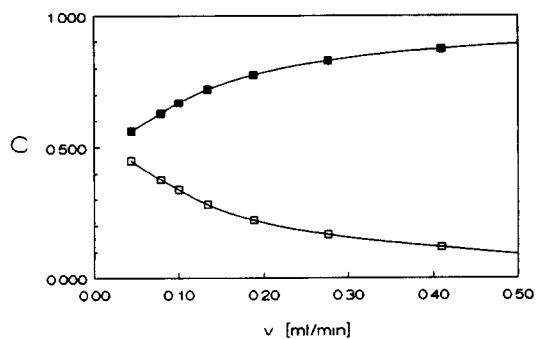


Fig 7 Steady-state concentration versus flow-rate in ml min^{-1} measured downstream of the (■) donor and (□) acceptor channels

ability was calculated by interpolation to be $(4.66 \pm 0.64) \times 10^{-6} \text{ m s}^{-1}$, which is at least of the same order of magnitude as the values obtained in the present study. Because the details of the experimental procedure for determining the permeability by the manufacturer are lacking and a different electrolyte at a different concentration is used, it is impossible to assess the causes of the discrepancy.

Influence of system parameters on sensitivity and sample throughput

After having proved the validity of the model, it is possible to investigate by numerical simulations the influence of the main design and operational parameters of a flow-injection system with a membrane separation module on the sensitivity and sample throughput of analysis. In this way costly and time-consuming experiments could be avoided. For convenience, the parameters of the experimental flow system outlined above could be used in the simulations. This will make it possible in some instances to compare simulated with experimental results. As the empirical flow-rate dependence of the axial dispersion and mass-transfer coefficients in the channels of the dialysis module was determined in the range $0-1.0 \text{ ml min}^{-1}$ for the former and $0-0.5 \text{ ml min}^{-1}$ for the latter coefficients, the simulations were confined to the narrower flow range. In an earlier study, the influence of the main parameters of a parallel-plate dialyser under laminar flow conditions was thoroughly investigated and conclusions for choosing their optimum values with respect to mass-transfer efficiency were drawn [21]. The calculations performed for a pulse concentration input were made under the assumption of zero length of the reactor and no dispersion in the injection section. Neither of these conditions holds in real flow-injection manifolds. For this reason, it is expedient to investigate how the dispersion processes in the flow sections upstream of the donor channel (i.e., fore-section, injection section and reactor) and the sample volume ($-X_e$) will influence the peak height (c_{max} or C_{max}) and the peak broadening (e or E) of the concentration profile monitored at the outlet of the acceptor channel. For flow systems with rela-

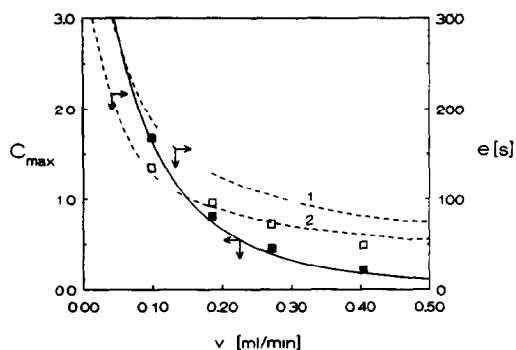


Fig 8 Flow-rate dependence of the maximum concentration (C_{max}) in the acceptor channel (solid line) and the peak width (dotted lines) at (1) $0.01 C_{\text{max}}$ and (2) $0.1 C_{\text{max}}$. Experimental results for (■) C_{max} and (□) for the peak width at $0.1 C_{\text{max}}$.

tively small sample volumes, the dispersion of the analyte in the reactor is the process playing a major role in the formation of the concentration profile at the inlet of the donor channel [17]. This process can be characterized by the reactor Fourier number (τ_r)

$$\tau_r = \pi D_r x_1 / v_D \quad (28)$$

From Eqn 28, it follows that concentration profile at the inlet of the donor channel for a given analyte can be easily manipulated by varying the volumetric flow-rate or (and) the length of the reactor. The volumetric flow-rate in the acceptor channel was assumed to be equal to that in the donor line (as in the experimental manifold). The flow-rate v_D is an operational parameter affecting considerably the dispersion processes in all sections of the flow system (including the channels of the dialysis module) and the rate of mass transfer. Its effect on peak height and peak broadening is illustrated in Fig 8. The peak broadening (e) is given in real dimensions (i.e., in seconds) to avoid confusion originating from the fact that on changing the flow-rate the mean residence time, which is the dimensionless time unit, also changes. The flow-rate dependences of C_{max} and e (at both $0.01 C_{\text{max}}$ and $0.1 C_{\text{max}}$) are similar, showing that in choosing the flow-rate a compromise should be made between sensitivity and sample throughput. The comparison of the theoretically calculated relationships for the flow-rate dependence of the maximum C_A and

the peak width (at C_A equal to 10% of the maximum C_A) with experimental results shows a satisfactory agreement (Fig 8) The agreement for e measured at 1% of the maximum C_A was poor because the signal monitored for such low concentrations (10^{-6} – 10^{-5} M) was in the range of the noise of the conductivity meter used

The influence of the reactor length (x_1) on c_{\max} and e (at 1% of c_{\max}) at a constant flow-rate (0.2 ml min^{-1}) is illustrated in Fig 9 These results show that shorter reactors are more favourable for the optimum performance of the flow-injection manifolds considered in this study However, it should be taken into consideration that a longer reactor can improve to a certain extent the reproducibility of analysis The reason for this is that the greater the extent to which the dispersion process is diffusion controlled (i.e., high τ), the smaller is the influence of arbitrary disturbances in the flow (e.g., slight variation in the flow-rate, pulsations) on the detected signal

The influence of the sample volume ($-X_e$) on the acceptor peak maximum ($c = -c_0/\gamma_1 X_e C_{\max}$, where c_0/γ_1 is constant throughout the simulations) and on the peak width (at 1% of c_{\max}) is presented in Fig 10 For simplicity in the calculations, the axial dispersion coefficient in the injection section was assumed to be equal to that of the reactor even for relatively large sample volumes (i.e., $-X_e > 0.1$) As expected, an increase in the sample volume leads to a higher peak maximum and hence to a higher sensitivity of analysis The dependence is almost linear for

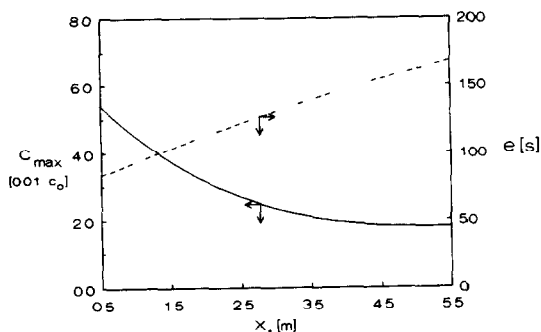


Fig 9 Influence of the reactor length (x_1) on c_{\max} (solid line) and e at $0.01c_{\max}$ (dotted line) $v_D = 0.2 \text{ ml min}^{-1}$

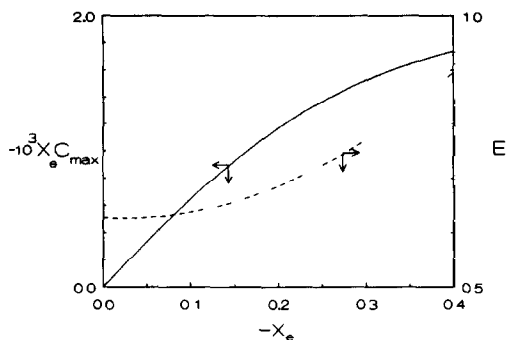


Fig 10 Influence of the sample volume ($-X_e$) on c_{\max} (solid line) and e at $0.01c_{\max}$ (dotted line) $V_D = 0.2 \text{ ml min}^{-1}$

dimensionless sample volumes up to 0.2, above which the slope gradually decreases and C_A approaches asymptotically the steady-state acceptor concentration for volumes higher than 0.8 [21] In the linear range mentioned above ($-X_e < 0.2$), the sample volume has virtually no influence on the sample throughput of the manifold For larger sample volumes the baseline-to-baseline time increases with $-X_e$ These results show that choosing a sample volume in the upper part of the linear range (e.g., 0.1–0.2) seems to be appropriate with respect to both sensitivity and sample throughput

Conclusions

A general mathematical model for a flow-injection system with a dialysis module based on the axially dispersed plug flow model was developed It was solved by Laplace transforms for values of the reactor Peclet number of less than 120, whereas for higher values a numerical method using the implicit alternating-direction finite difference method was found to be more suitable

The validity of the model was confirmed experimentally for a flow-injection manifold utilizing a parallel-plate dialysis model This geometry of the dialyser does not limit the generality of the model, which can also be applied to other geometries of the membrane separation module (e.g., tubular) and other types of membranes (e.g., gas-diffusion)

The results obtained in this study show that the proposed model for a flow-injection manifold with a membrane separation module, together with the parameter identification procedure outlined above, can be used successfully for a better understanding of the processes taking place in such manifolds and for predicting the shape and magnitude of the analytical signal, for optimizing the performance of these manifolds with respect to sensitivity and sample throughput, for characterizing various membranes (i.e., determining their thickness and the diffusivity of different analytes in them) and for characterizing the mass-transfer properties of various geometries of the donor and the acceptor channels (i.e., determining their axial dispersion and mass-transfer coefficients). The last two possible applications of the model could be useful in membrane technology and in designing effective process, medical (for haemodialysis) or analytical flow-through dialysers.

The possibility provided by simulating the

model to select the values of some construction and operational parameters of the flow-injection systems incorporating membrane separation modules which will ensure their optimum performance is illustrated on the experimental manifold used for verification of the model. On the basis of the simulation results, some more general conclusions concerning the selection of optimum sample volume and reactor length were drawn: the sample volume ($-X_e$) should be in the range 0.1-0.3 and the reactor length should not exceed the value necessary for obtaining a reproducible concentration profile at the inlet of the donor channel. The last condition is usually in effect for diffusion-controlled dispersion ($\tau_r > 0.7$). The guidelines for choosing optimum values for the parameters of the membrane separation module were formulated elsewhere [21].

The authors are grateful to Egbert Hoogendam for experimental help.

APPENDIX A

Laplace domain solution of Eqns 3-11

The Laplace transforms of Eqns 3-11 are the following

For the donor line

$$\frac{d^2 \bar{C}_f}{dX^2} - P_f \frac{d\bar{C}_f}{dX} - \frac{pP_f}{\gamma_f} \bar{C}_f = 0 \quad X \leq X_e \quad (\text{A1})$$

$$\frac{d^2 \bar{C}_1}{dX^2} - P_1 \frac{d\bar{C}_1}{dX} - \frac{pP_1}{\gamma_1} \bar{C}_1 = -\frac{P_1}{\gamma_1} \quad X_e \leq X \leq 0 \quad (\text{A2})$$

$$\frac{d^2 \bar{C}_r}{dX^2} - P_r \frac{d\bar{C}_r}{dX} - \frac{pP_r}{\gamma_r} \bar{C}_r = 0 \quad 0 \leq X \leq X_1 \quad (\text{A3})$$

$$\frac{d^2 \bar{C}_D}{dX^2} - P_D \frac{d\bar{C}_D}{dX} - (p + K_D) \frac{P_D}{\gamma_D} \bar{C}_D + K_D \frac{P_D}{\gamma_D} \bar{C}_{Dm} = 0 \quad X_1 \leq X \leq X_2 \quad (\text{A4})$$

$$\frac{d^2 \bar{C}_a}{dX^2} - P_a \frac{d\bar{C}_a}{dX} - \frac{pP_a}{\gamma_a} \bar{C}_a = 0 \quad X \geq X_2 \quad (\text{A5})$$

for the membrane

$$\frac{d^2 \bar{C}_m}{dY^2} - P_m p \bar{C}_m = 0 \quad 2 \leq Y \leq 2 + 2\Delta \quad (\text{A6})$$

For the acceptor line

$$\frac{d^2 \bar{C}_f^*}{dX^2} - P_f^* \frac{d\bar{C}_f^*}{dX} - \frac{pP_f^*}{\gamma_f^*} \bar{C}_f^* = 0 \quad X \leq X_1 \quad (\text{A7})$$

$$\frac{d^2 \bar{C}_A}{dX^2} - P_A \frac{d\bar{C}_A}{dX} - (p + K_A) \frac{P_A}{\gamma_A} \bar{C}_A + K_A \frac{P_A}{\gamma_A} \bar{C}_{Am} = 0 \quad X_1 \leq X \leq X_2 \quad (\text{A8})$$

$$\frac{d^2 \bar{C}_a^*}{dX^2} - P_a^* \frac{d\bar{C}_a^*}{dX} - \frac{pP_a^*}{\gamma_a^*} \bar{C}_a^* = 0 \quad X \geq X_2 \quad (\text{A9})$$

The corresponding boundary conditions can be obtained from those of Eqns 3-11 by simply replacing the concentrations with their Laplace transforms

The Laplace domain solutions of Eqns A1-A9 are the following

$$\bar{C}_f = A_1 \exp[P_f(0.5 + q_f)(X - X_e)] \quad (\text{A10})$$

$$\bar{C}_i = A_2 \exp[P_i(0.5 + q_i)X] + A_3 \exp[P_i(0.5 - q_i)X] + 1/p \quad (\text{A11})$$

$$\bar{C}_r = A_4 \exp[P_r(0.5 + q_r)X] + A_5 \exp[P_r(0.5 - q_r)X] \quad (\text{A12})$$

$$\bar{C}_D = \sum_{i=1}^{i=4} r_i \exp[z_i(X - X_1)] \quad (\text{A13})$$

$$\bar{C}_a = A_6 \exp[P_a(0.5 - q_a)(X - X_2)] \quad (\text{A14})$$

$$\bar{C}_f^* = B_1 \exp[P_f^*(0.5 + q_f^*)X] \quad (\text{A15})$$

$$\bar{C}_A = \sum_{i=1}^{i=4} \alpha_i r_i \exp[z_i(X - X_1)] \quad (\text{A16})$$

$$\bar{C}_a^* = B_2 \exp[P_a^*(0.5 - q_a^*)(X - X_2)] \quad (\text{A17})$$

where

$$q_j = \left(\frac{p}{\gamma_j P_j} + 0.25 \right)^{1/2} \quad (\text{A18})$$

The coefficients A_1 - A_6 , B_1 , B_2 and r_i are the unknown integrational constants which can be obtained from the boundary conditions of Eqns A1-A5 and A7-A9. The coefficients α_i can be determined from solution of Eqn A6 [21]

$$\alpha_i = - \frac{z_i^2 - P_D z_i - [p + K_D(1 - f_D)] \frac{P_D}{\gamma_D}}{f_A K_D \frac{P_D}{\gamma_D}} \quad (\text{A19})$$

where z_i are the roots of the following biquadratic equation [21]

$$z^4 - (P_D + P_A)z^3 + \left(P_D P_A - [p + K_D(1 - f_D)] \frac{P_D}{\gamma_D} - [p + K_A(1 - g_A)] \frac{P_A}{\gamma_A} \right) z^2 + P_D P_A \left(\frac{p + K_D(1 - f_D)}{\gamma_D} + \frac{p + K_A(1 - g_A)}{\gamma_A} \right) z + \frac{P_D}{\gamma_D} \frac{P_A}{\gamma_A} ([p + K_D(1 - f_D)][p + K_A(1 - g_A)] - K_D K_A f_A g_D) = 0 \tag{A20}$$

In the general case when the overall mass-transfer process in the membrane separation module is governed by the transfer both in the channels and in the membrane, the coefficients f_D , f_A , g_D and g_A are defined as follows [21]

$$f_D = \frac{H_D}{F_D Q} \{ (q_m + H_A/F_A) \exp(2q_m \Delta) + (q_m - H_A/F_A) \exp(-2q_m \Delta) \}$$

$$f_A = (H_A/F_D Q) 2q_m$$

$$g_D = (H_D/F_A Q) 2q_m$$

$$g_A = (H_A/F_A Q) [(q_m + H_D/F_D) \exp(2q_m \Delta) + (q_m - H_D/F_D) \exp(-2q_m \Delta)]$$

where

$$Q = (q_m + H_D/F_D)(q_m + H_A/F_A) \exp(2q_m \Delta) - (q_m - H_D/F_D)(q_m - H_A/F_A) \exp(-2q_m \Delta)$$

$$q_m = (pP_m)^{1/2}$$

When the overall mass transfer is controlled only by the transfer in the channels of the membrane separation module, e.g., the resistance of the membrane is negligible because $\Delta \rightarrow 0$, the expressions for f_D , f_A , g_D and g_A can be considerably simplified

$$f_D = \left(1 + \frac{F_D}{F_A} \frac{K_A}{k_D} \right)^{-1}$$

$$f_A = \left(\frac{F_D}{F_A} + \frac{k_D}{k_A} \right)^{-1}$$

$$g_D = \left(\frac{F_A}{F_D} + \frac{k_A}{k_D} \right)^{-1}$$

$$g_A = \left(1 + \frac{F_A}{F_D} \frac{k_D}{k_A} \right)^{-1}$$

Another extreme case is when the resistance of the membrane is much higher than that of the channels of the membrane separation module. It can be assumed that the mass-transfer coefficients of both channels are infinitely high, i.e., $k_D \rightarrow \infty$ and $k_A \rightarrow \infty$. The following relationships necessary for the solution of Eqn. A20 will hold

$$K_D f_A = F_A q_m / [2P_m \sinh(2q_m \Delta)]$$

$$K_D (1 - f_D) = F_D q_m / [2P_m \tanh(2q_m \Delta)]$$

$$K_A g_D = (a_D/a_A) F_D q_m / [2P_m \sinh(2q_m \Delta)]$$

$$K_A (1 - g_A) = (a_D/a_A) F_A q_m / [2P_m \tanh(2q_m \Delta)]$$

The Laplace transforms of the concentration profiles in the flow sections situated upstream (i.e., reactor) and downstream (i.e., after-section of the donor line) of the donor channel and downstream of the acceptor channel (i.e., after-section of the acceptor line) are the following

$$\bar{C}_r = \frac{1}{2q_r} \left\{ \left[\sum_{i=1}^{i=4} r_i (0.5 + q_r - z_i/P_D) \right] \exp[P_r(0.5 - q_r)(X - X_1)] - \left[\sum_{i=1}^{i=4} r_i (0.5 - q_r - z_i/P_D) \right] \exp[P_r(0.5 + q_r)(X - X_1)] \right\} \quad 0 \leq X \leq X_1 \quad (\text{A21})$$

$$\bar{C}_a = \left\{ \sum_{i=1}^{i=4} r_i \exp[z_i(X_2 - X_1)] \right\} \exp[P_a(0.5 - q_a)(X - X_2)] \quad X \geq X_2 \quad (\text{A22})$$

$$\bar{C}_a^* = \left\{ \sum_{i=1}^{i=4} \alpha_i r_i \exp[z_i(X_2 - X_1)] \right\} \exp[P_a^*(0.5 - q_a^*)(X - X_2)] \quad X \geq X_2 \quad (\text{A23})$$

where

$$\begin{aligned} r_1 &= h - r_2 h_1^1 - r_3 h_2^1 - r_4 h_3^1 \\ r_2 &= \frac{h + r_3(h_2^2 - h_2^1) + r_4(h_3^2 h_3^1)}{h_1^1 - h_1^2} \\ r_3 &= \frac{h(h_1^2 - h_1^1) + r_4[(h_1^2 - h_1^1)(h_3^4 - h_3^1) - (h_1^4 - h_1^1)(h_3^2 - h_3^1)]}{(h_1^4 - h_1^1)(h_2^2 - h_2^1) - (h_1^2 - h_1^1)(h_2^4 - h_2^1)} \\ r_4 &= \left[h[(h_1^2 - h_1^1)[(h_3^2 - h_2^1)(h_1^2 - h_1^1) - (h_3^4 - h_3^1)(h_2^2 - h_2^1)] - (h_1^2 - h_1^1)[(h_1^2 - h_1^1)(h_2^4 - h_2^1) - (h_2^2 - h_2^1)(h_1^4 - h_1^1)] \right] \\ &\quad \times \left[[(h_3^3 - h_3^1)(h_1^2 - h_1^1) - (h_3^2 - h_3^1)(h_1^3 - h_1^1)] [(h_1^2 - h_1^1)(h_2^4 - h_2^1) - (h_2^2 - h_2^1)(h_1^4 - h_1^1)] - [(h_1^2 - h_1^1)(h_3^4 - h_3^1) - (h_2^2 - h_2^1)(h_1^3 - h_1^1)] [(h_1^2 - h_1^1)(h_2^3 - h_2^1) - (h_2^2 - h_2^1)(h_1^3 - h_1^1)] \right]^{-1} \end{aligned}$$

$$h = 2q_r M_3 \exp[P_r(0.5 + q_r)X_1] / S_1$$

$$h_j^1 = S_{j+1} / S_1$$

$$h_j^2 = [(0.5 - q_a - z_{j+1}/P_D) / (0.5 - q_a - z_1/P_D)] \exp[(z_{j+1} - z_1)(X_2 - X_1)]$$

$$h_j^3 = [\alpha_{j+1}(0.5 + q_f^* - z_{j+1}/P_A)] / [\alpha_1(0.5 + q_f^* - z_1/P_A)]$$

$$h_j^4 = [\alpha_{j+1}(0.5 - q_a^* - z_{j+1}/P_A)] / [\alpha_1(0.5 - q_a^* - z_1/P_A)] \exp[(z_{j+1} - z_1)(X_2 - X_1)]$$

where $j = 1, 2, 3$,

$$S_j = 2q_r M_1 + [M_2 \exp[2P_r q_r X_1] - M_1](0.5 + q_r - z_j/P_D)$$

where $j = 1, 2, 3, 4$,

$$M_1 = (q_r - q_i)(q_f + q_i) \exp[-2P_r q_i X_e] - (q_r + q_i)(q_f - q_i)$$

$$M_2 = (q_r - q_i)(q_f - q_i) - (q_f + q_i)(q_r + q_i) \exp[-2P_r q_i X_e]$$

$$M_3 = [(0.5 - q_i)(q_f - q_i) + 2q_i(0.5 + q_f) \exp[-P_r(0.5 + q_i)X_e] - (0.5 + q_i)(q_f + q_i) \exp[-2P_r q_i X_e]] / p$$

If the membrane is impermeable to the analyte, the dispersion process in the donor channel will be described by Eqn A24 instead of Eqn A13

$$\bar{C}_D = A_7 \exp[P_D(0.5 + q_D)(X - X_1)] + A_8 \exp[P_D(0.5 - q_D)(X - X_1)] \tag{A24}$$

The Laplace transform of the concentration profile monitored in the donor line downstream of the donor channel (i.e., in the after-section) will be given by Eqn A14, where

$$A_6 = \frac{1}{p} \{ 4q_r q_D \exp[P_r(0.5 - q_r)X_1 + P_D(0.5 - q_D)(X_2 - X_1)] \\ \times \{ (0.5 + q_1)(q_f + q_1) - (0.5 - q_1)(q_f - q_1) \exp[2P_1 q_1 X_e] - 2q_1(0.5 + q_f) \exp[-P_1(0.5 - q_1)X_e] \} \\ \times \{ (q_a + q_D)(q_r + q_D) - (q_a - q_D)(q_r - q_D) \exp[-2P_D q_D (X_2 - X_1)] \} \\ \times \{ (q_r + q_1)(q_f + q_1) - (q_f - q_1)(q_r - q_1) \exp[-2P_1 q_1 X_e] \} - \{ (q_r - q_1)(q_f + q_1) - (q_r + q_1) \\ \times (q_f - q_1) \exp[-2P_1 q_1 X_e] \} \\ \times \{ (q_a + q_D)(q_r - q_D) - (q_a - q_D)(q_r + q_D) \exp[-2P_D q_D (X_2 - X_1)] \} \exp[-2P_r q_r X_1] \}^{-1} \tag{A25}$$

APPENDIX B

Implicit finite-difference equations for the membrane separation module

The θ , X and Y regions of interest were divided into K , N and M sub-intervals, respectively, so that it can be written

$$\theta = k \Delta\theta \quad (k = 0, 1, 2, \dots, K)$$

$$X = i \Delta X \quad (i = 0, 1, 2, \dots, N)$$

$$Y = j \Delta Y \quad (j = 0, 1, 2, \dots, M)$$

X-implicit equations

$$-\frac{\gamma_D \Delta\theta}{4 \Delta X} \left(1 + \frac{1}{P_D \Delta X} \right) C_{D_{i-1k+0.5}} + \left[1 + \frac{\Delta\theta \gamma_D}{(\Delta X)^2 P_D} \right] C_{D_{ik+0.5}} + \frac{\gamma_D \Delta\theta}{4 \Delta X} \left(1 - \frac{2}{P_D \Delta X} \right) C_{D_{i+1k+0.5}} \\ = \left(1 - \frac{K_D \Delta\theta}{2} \right) C_{D_{ik}} + \frac{K_D \Delta\theta}{2F_D} C_{m_{0k}} \tag{B1}$$

$$-\frac{\gamma_A \Delta\theta}{4 \Delta X} \left(1 + \frac{2}{P_A \Delta X} \right) C_{A_{i-1k+0.5}} + \left[1 + \frac{\Delta\theta \gamma_A}{(\Delta X)^2 P_A} \right] C_{A_{ik+0.5}} + \frac{\gamma_A \Delta\theta}{4 \Delta X} \left(1 - \frac{2}{P_A \Delta X} \right) C_{A_{i+1k+0.5}} \\ = \left(1 - \frac{K_A \Delta\theta}{2} \right) C_{A_{ik}} + \frac{K_A \Delta\theta}{2F_A} C_{m_{Mk}} \tag{B2}$$

$$C_{m_{0k+0.5}} = \frac{1}{1 + H_D \Delta Y / F_D} (C_{m_{1k+0.5}} + H_D \Delta Y C_{D_{ik+0.5}}) \quad \text{for } j = 0 \tag{B3}$$

$$C_{m_{j,k+0.5}} = \left(1 - \frac{\Delta\theta}{(\Delta Y)^2 P_m} \right) C_{m_{j,k}} + \frac{\Delta\theta}{2(\Delta Y)^2 P_m} (C_{m_{j+1,k}} + C_{m_{j-1,k}}) \quad \text{for } 0 < j < M \tag{B4}$$

$$C_{m_{Mk+0.5}} = \frac{1}{1 + H_A \Delta Y / F_A} (C_{m_{M-1k+0.5}} + H_A \Delta Y C_{A_{ik+0.5}}) \quad \text{for } j = M \tag{B5}$$

Y-implicit equations

$$C_{D_i, k+1} = \frac{\gamma_D \Delta \theta}{4 \Delta X} \left(1 + \frac{2}{P_D \Delta X} \right) C_{D_{i-1}, k+0.5} + \left[1 - \frac{\Delta \theta \gamma_D}{(\Delta X)^2 P_D} - \frac{K_D \Delta \theta}{2} \right] C_{D_i, k+0.5} - \frac{\gamma_D \Delta \theta}{4 \Delta X} \left(1 - \frac{2}{P_D \Delta X} \right) C_{D_{i+1}, k+0.5} + \frac{K_D \Delta \theta}{2 F_D} C_{m, 0, k+0.5} \quad (\text{B6})$$

$$C_{A_i, k+1} = \frac{\gamma_A \Delta \theta}{4 \Delta X} \left(1 + \frac{2}{P_A \Delta X} \right) C_{A_{i-1}, k+0.5} + \left[1 - \frac{\Delta \theta \gamma_A}{(\Delta X)^2 P_A} - \frac{K_A \Delta \theta}{2} \right] C_{A_i, k+0.5} - \frac{\gamma_A \Delta \theta}{4 \Delta X} \left(1 - \frac{2}{P_A \Delta X} \right) C_{A_{i+1}, k+0.5} + \frac{K_A \Delta \theta}{2 F_A} C_{m, M, k+0.5} \quad (\text{B7})$$

$$\left(1 + \frac{H_D \Delta Y}{F_D} \right) C_{m, 0, k+1} - C_{m, 1, k+1} = H_D \Delta Y C_{D_i, k+1} \quad \text{for } j = 0 \quad (\text{B8})$$

$$\frac{\Delta \theta}{2(\Delta Y)^2 P_m} C_{m, j-1, k+1} - \left[1 + \frac{\Delta \theta}{(\Delta Y)^2 P_m} \right] C_{m, j, k+1} + \frac{\Delta \theta}{2(\Delta Y)^2 P_m} C_{m, j+1, k+1} = -C_{m, j, k+0.5} \quad \text{for } 0 < j < M \quad (\text{B9})$$

$$-C_{m, M-1, k+1} + \left(1 + \frac{H_A \Delta Y}{F_A} \right) C_{m, M, k+1} = H_A \Delta Y C_{A_i, k+1} \quad \text{for } j = M \quad (\text{B10})$$

Both the sets of *X*-implicit and *Y*-implicit equations have a tridiagonal coefficient matrix and were solved in a straightforward way by the Gaussian elimination method [25]

APPENDIX C

Equations for calculating the steady-state membrane permeability

The concentrations at the outlets of the donor and the acceptor channels under steady-state conditions (i.e., $\partial C / \partial \theta = 0$) can be calculated by the following equations [29]

$$C_D = \frac{\omega}{\omega - F_D / F_A} \left\{ 1 - \frac{F_D}{\omega F_A} \exp[\beta(X_2 - X_1)] \right\} \quad (\text{C1})$$

$$C_A = \frac{\omega}{\omega F_A / F_D - 1} \left\{ 1 - \exp[\beta(X_2 - X_1)] \right\} \quad (\text{C2})$$

The coefficient β is the only negative root of Eqn C3 and can be calculated analytically by Cardan's equation [30]

$$z^3 - (P_D + P_A)z^2 + \left[P_D P_A - \Psi \left(\frac{v_A}{v_D} F_D P_D + F_A P_A \right) \right] z + \Psi P_D P_A \left(\frac{v_A}{v_D} F_D + F_A \right) = 0 \quad (\text{C3})$$

where

$$\Psi = \left[\frac{v_A}{v_D} \gamma_D \left(\frac{a_D}{a_A} \frac{F_A}{K_A} + \frac{F_D}{K_D} + 4P_m \Delta \right) \right]^{-1} \quad (C4)$$

Once determined, β can be used for calculating the coefficient ω in Eqns C1 and C2

$$\omega = \frac{v_D \beta}{v_A F_A \Psi} \left(1 - \frac{\beta}{P_D} \right) + \frac{F_D}{F_A} \quad (C5)$$

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